

## Infrared Spectra and Molecular Parameters of Matrix-isolated Gallium(I), Indium(I), and Thallium(I) Oxides (Ga<sub>2</sub>O, In<sub>2</sub>O, and Tl<sub>2</sub>O)

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*Summary* The  $\nu_3$  fundamentals are observed for <sup>16</sup>O and <sup>18</sup>O species, and these frequency ratios used directly to derive bond angles: principal stretching force constants are evaluated assuming a range of interaction constants, and the corresponding bond lengths estimated *via* the Laurie-Herschbach relationship.

In recent years the combined techniques of matrix isolation and i.r. spectroscopy have provided valuable structural information on high-temperature chemical species.<sup>1</sup> We describe part of an investigation into the vaporisation and condensation of Group III oxides. It is known from mass-spectrometric studies<sup>2</sup> that the suboxides M<sub>2</sub>O (M = Ga, In, Tl) are present in the vapour above samples of M<sub>2</sub>O<sub>3</sub>. However, apart from a brief communication describing the low-resolution vapour-phase i.r. spectra<sup>3</sup> of Ga<sub>2</sub>O and In<sub>2</sub>O, and an electron-diffraction study<sup>4</sup> of Ga<sub>2</sub>O and In<sub>2</sub>O, the interpretation of which has recently been questioned,<sup>5</sup> little is known about these species, and no force-constant data are available.

When the vapour species present above heated Ga<sub>2</sub>O<sub>3</sub> are trapped in a nitrogen matrix at 15°K, the i.r. spectrum shows an intense band at 809.4 cm.<sup>-1</sup>. <sup>18</sup>O-enrichment gives rise to only one additional band at 768.5 cm.<sup>-1</sup> and the relative intensity of these two bands reflects the original <sup>16</sup>O:<sup>18</sup>O isotopic distribution. This molecular species evidently contains only one atom of oxygen. It

cannot be diatomic GaO, since Ga<sup>16</sup>O is known<sup>6</sup> to absorb at *ca.* 755 cm.<sup>-1</sup>, and on the basis of the mass-spectrometric data,<sup>2</sup> can only be Ga<sub>2</sub>O. Similar results were obtained for the In/O and Tl/O systems, and the frequencies observed are listed in the Table. A number of very weak bands were also present in some spectra: these will be discussed in a later paper. The same results were obtained when oxygen gas was passed over the heated metals, and this proved the most satisfactory method of obtaining <sup>18</sup>O enrichment.

Matrix-isolation studies on the Al<sub>2</sub>O<sub>3</sub>/Al system<sup>7</sup> have shown that in the i.r. spectrum of Al<sub>2</sub>O, the strongest band is the antisymmetric stretch  $\nu_3$ . The bands observed in this work are therefore assigned as being the antisymmetric modes of Ga<sub>2</sub>O, In<sub>2</sub>O, and Tl<sub>2</sub>O. For triatomic molecules M<sub>2</sub>O with symmetries C<sub>2v</sub> or D<sub>∞h</sub>, the zero-order antisymmetric vibration frequency  $\omega_3$  is given by:

$$4\pi^2 \omega_3^2 = (F_r - F_{rr})(\mu_M + 2\mu_O \sin^2\theta) \quad (1)$$

where  $\mu_M$  and  $\mu_O$  are the reciprocal masses of atoms M and O, and  $2\theta$  is the angle M-O-M.  $F_r$  is the principal stretching force constant, and  $F_{rr}$  is an interaction constant. From equation (1) it is evident that:

$$\left(\frac{\omega_3}{\omega_3^i}\right)^2 = \frac{[\mu_M + 2\mu(^{16}\text{O}) \sin^2\theta]}{[\mu_M + 2\mu(^{18}\text{O}) \sin^2\theta]} \quad (2)$$

where  $\omega_3^i$  refers to the isotopically substituted M<sub>2</sub><sup>18</sup>O.

Although this equation holds rigorously only for zero-order frequencies, it may usefully be applied to estimate a value for  $\theta$  from observed i.r. transitions provided that anharmonicity corrections and matrix shifts are small.<sup>7,8</sup>

The M-O-M angles given in the Table are calculated

be used to estimate the bond length  $r_{M-O}$  via the appropriate Laurie-Herschbach relationship.<sup>9</sup> The bond lengths listed in the Table have been derived in this way, and correspond closely to the values obtained from electron diffraction.

TABLE  
I.r. absorptions and molecular parameters of Ga<sub>2</sub>O, In<sub>2</sub>O, and Tl<sub>2</sub>O

		$\nu_3 M_2^{16}O$ (cm. <sup>-1</sup> )	$\nu_3 M_2^{18}O$ (cm. <sup>-1</sup> )	MOM	$F_r$ (mdynes/Å)	$r_{M-O}$ (Å)
M = Ga	This work	809.4 ± 0.1	768.5 ± 0.1	143 ± 5°	3.05 ± 0.35	1.86 ± 0.03
	Previous work	770 <sup>a</sup>	—	140 ± 10° <sup>b</sup> 150° <sup>c</sup>	—	1.84 ± 0.01 <sup>b</sup>
M = In	This work	722.4 ± 0.1	684.3 ± 0.1	135 ± 7°	2.67 ± 0.3	2.04 ± 0.04
	Previous work	680 <sup>a</sup>	—	150 ± 10° <sup>b</sup>	—	2.02 ± 0.01 <sup>b</sup>
M = Tl	This work	625.3 ± 0.1	591.2 ± 0.1	131 ± 11°	2.12 ± 0.25	2.19 ± 0.05
	Previous work	—	—	130° <sup>c</sup>	—	—

<sup>a</sup> ref. 3    <sup>b</sup> ref. 4    <sup>c</sup> ref. 5

directly from the observed frequencies using equation (2). They correlate well with a statement made by Rambidi *et al.*<sup>5</sup> in a re-assessment of the original electron-diffraction data,<sup>4</sup> although the observed trend may not be significant in view of the approximations involved. The principal force constant  $F_r$  may be obtained directly from equation (1) by setting  $F_{rr} = \pm 0.1F_r$ . This parameter may then

These results indicate that all three fundamental vibrations of Ga<sub>2</sub>O, In<sub>2</sub>O, and Tl<sub>2</sub>O are i.r.-active. A knowledge of these frequencies will provide important thermodynamic data for these high-temperature species and will also allow a complete determination of the quadratic force field.

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